Short-term and spike-timing-dependent plasticities facilitate the formation of modular neural networks

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The brain has the phenomenal ability to reorganize itself by forming new connections among neurons and by pruning others. The so-called neural or brain plasticity facilitates the modification of brain structure and function over different time scales. Plasticity might occur due to external stimuli received from the environment, during recovery from brain injury, or due to modifications within the body and brain itself. In this paper, we study the combined effect of short-term (STP) and spike-timing-dependent plasticities (STDP) on the synaptic strength of excitatory coupled Hodgkin-Huxley neurons and show that plasticity can facilitate the formation of modular neural networks with complex topologies that resemble those of networks with preferential attachment properties. In particular, we use an STDP rule that alters the synaptic coupling intensity based on time intervals between spikes of postsynaptic and presynaptic neurons. Previous works have shown that STDP may induce the appearance of directed connections from high to low frequency spiking neurons. On the other hand, STP is attributed to the release of neurotransmitters in the synaptic cleft of neurons that alter its synaptic efficiency. Our results suggest that the combined effect of STP and STDP with high recovery time facilitates the formation of connections among neurons with similar spike frequencies only, a kind of preferential attachment. We then pursue this further and show that, when starting with all-to-all neural configurations, depending on the STP recovery time and distribution of neural frequencies, modular neural networks can emerge as a direct result of the combined effect of STP and STDP.

Neural or brain plasticity is the remarkable ability of the brain to alter its structure and function over time to achieve cognitive functions and perform tasks. Even though, plasticity may result as a consequence of external stimuli, it might also occur during learning processes and after brain damage. Synaptic plasticity refers to the ability of the brain to render synapses weaker or stronger according to neural activity. In this paper, we build a neural network model of excitatory coupled Hodgkin-Huxley (HH) neurons with the ability to modify their synaptic strengths over time, i.e., a neural network with short-term (STP) and spike-timing-dependent plasticities (STDP) build in. STDP acts on longer time scales compared to STP, with both plasticities playing an important role in brain functions. Here, we show that due to STP, neural networks equipped

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with STDP facilitate the formation of synapses among neurons with similar spike frequencies only, a kind of preferential attachement. Modular neural networks can emerge as a direct result of the combined effect of STP and STDP, a structure also depicted by real brain networks.

I. INTRODUCTION

Mathematical models have been employed in neuroscience since early in the 20th century¹ to explain experimental findings and elucidate the inner workings of the brain. In 1907, Lapicque² proposed the integrate-and-fire model that can reproduce the electrical activity of the membrane potential of neurons. It is one of the most popular models for studying the dynamic behaviour of neural systems. Later in 1952, Hodgkin and Huxley³ explained the ionic mechanisms in the cell membrane of neurons and proposed the so-called Hodgkin-Huxley (HH) neural model that has, since then, been used extensively to study neural networks⁴.

In particular, neural network models have been used extensively in neuroscience applications, such as in studies of neural dynamics⁵, dynamic range⁶⁻⁹, neural synchronization¹⁰⁻¹⁴, flow of information¹⁵⁻¹⁷ and brain plasticity^{18,19} to name a few.

Neural plasticity is the ability of the brain to modify its function and structure over different time scales^{21,22}. The term was initially used by James²³ in 1890 to propose that phenomena of habit in living organisms are due to plasticity. Cajal²⁴ reported in the early 1900s his research about regenerative and degeneration changes in brain structure. In 1924, Lashley²⁵ demonstrated experimental evidence of a malleable brain. Konorski²⁶ and Hebb²⁷ proposed in 1948 and 1949, respectively, that neural activities have influence on the connection among neurons. In 1960, Bennet et al.²⁸ carried out experiments with rats and observed chemical and anatomical plasticity in the brain. Since then, there have been many theoretical²⁹ and empirical³⁰ studies aiming to explain and understand brain plasticity and its effects in brain structure and function.

In synaptic plasticity, the synapses among neurons are potentiated or depressed in time according to the activity of the neurons^{31,32}. Recent works^{33,34} have shown that short-term (STP) and spike-timing-dependent plasticities (STDP) are different forms of neural processes leading to synaptic modifications. In particular, STDP depends on the relative timing of presynaptic and postsynaptic neural spikes^{35,36}. This type of plasticity can lead to various dynamical phenomena and coupling structures, such as stable localized structures³⁷, stimulation-induced synchronization or desynchronization³⁸, noise-enhanced synchronization^{39,40}, and nontrivial topology²⁰. The STDP mechanism plays a role in temporal coding of information by spikes^{37,41}. On the other hand, STP is attributed to the release of neurotransmitters in the synaptic cleft of neurons that alter its synaptic efficiency and acts on shorter time scales, ranging from milliseconds to hundreds to thousands of milliseconds⁴²⁻⁴⁵. As in the case with STDP, STP can have a great influence on the network's dynamical behavior. For instance, it may stabilize the parametric working memory⁴⁶, contribute to the emergence of spontaneous traveling waves⁴⁷, or induce phase changes in neural postsynaptic spiking³³.

Here, we extend the work in Borges et al.²⁰ which was focused on STDP only and study the combined effect of STP and STDP in neural networks of excitatory coupled HH neurons. The plasticity terms that model STDP in the equations in Sec. II are based on the experimental results by Bi and Poo^{48,49} that were performed on excitatory synapses and on theoretical results by Abbott et al.⁴³ and Popovych et al.³⁹ (STDP). The results in Bi and Poo^{48,49} show that STDP is a function of the relative timing of postsynaptic and presynaptic spikes and is theoretically backed by the Hebbian synaptic learning rule²⁷. Instead, STP depends on the neural recovery dynamics^{50,51}. McDonnell and Graham³³ used mathematical analysis and numerical simulations to show that STP induces phase changes in neural postsynaptic spiking. In our work, we start by studying the simplest case of a pair of HH neurons for a range of spike frequencies, aiming to understand how connectivity between them changes by the combined effect of STP and STDP. Next, we build an initially, all-to-all (globally) connected network of HH neurons and consider the simultaneous effect of STP and STDP in a range of coupling strengths. We show that STP plays an important role in topology changes in neural networks with STDP. Indeed, we find that for high STP recovery time, only neurons with similar spike frequencies tend to connect, a form of preferential attachement. More importantly, our results show that, when starting with all-to-all networks, depending on the STP recovery time and distribution of neural frequencies, modular neural networks can emerge as a direct result of the combined effect of STP and STDP, a structure depicted by neurophysiological and experimental studies^{52,53}. For the considered setup, STP plays a balancing role: while STDP tends to synchronize all neurons in one cluster, the STP destroys the strong synchronization and leads to a modular structure.

The paper is structured as follows: In Section II, we introduce the general mathematical model of HH neural networks with STP and STDP and, in Section III, we present our analysis and results based on numerical simulations, that show the effects of both plasticities, initially on a pair of neurons and then, on a network of 100 HH neurons. Finally, we present the conclusions of our work in the last section.

II. A HODGKIN-HUXLEY NEURAL NETWORK WITH STP AND STDP

We build use a neural network model of N HH neurons coupled with excitatory chemical synapses, equipped with STP and STDP rules based on the experimental results by Bi and Poo^{48,49} and theoretical models proposed by Popovych et al.³⁹ (STDP) and McDonnell and Graham^{33,43,50} (STP).

Specifically, the HH neural network model considered is given by

$$C\dot{V}_{i} = I_{i} - g_{\rm K} n_{i}^{4} (V_{i} - E_{\rm K}) - g_{\rm Na} m_{i}^{3} h_{i} (V_{i} - E_{\rm Na}) - g_{I} (V_{i} - E_{\rm I}) + (V_{r} - V_{i}) \sum_{i}^{N} \varepsilon_{ii} f_{i} D_{i},$$
(1)

$$-g_L(V_i - E_L) + (V_r - V_i) \sum_{j=1}^{\infty} \varepsilon_{ij} J_j D_j, \tag{1}$$

$$\dot{n}_{i} = \alpha_{n_{i}}(V_{i})(1 - n_{i}) - \beta_{n_{i}}(V_{i})n_{i}, \tag{2}$$

$$\dot{m}_i = \alpha_{m_i}(V_i)(1 - m_i) - \beta_{m_i}(V_i)m_i,$$
(3)

$$\dot{h}_i = \alpha_{h_i}(V_i)(1 - h_i) - \beta_{h_i}(V_i)h_i, \qquad (4)$$

$$\dot{f}_i = -\frac{f_i}{\tau},\tag{5}$$

$$\dot{D}_i = \frac{1 - D_i}{\tau_D},\tag{6}$$

where C (μ F/cm²) is the membrane capacitance and V_i (mV) the membrane potential of neuron *i* at time *t* (where i = 1, ..., N). I_i (μ A/cm²) is the constant current density of neuron *i* and ε_{ij} represents the matrix of coupling weights between neurons *i* and *j*. n_i and m_i are the activation of potassium and sodium functions, respectively, and h_i the inactivation of sodium function. Parameters *g* and *E* are associated with the conductance and reversal potential of each ion, respectively, and V_r is the excitatory reversal potential. The various rate functions in Eqs. (1) - (6) are given by

$$\alpha_n(v) = \frac{0.01v + 0.55}{1 - \exp\left(-0.1v - 5.5\right)},\tag{7}$$

$$\beta_n(v) = 0.125 \exp\left(\frac{-v - 65}{80}\right),\tag{8}$$

$$\alpha_m(v) = \frac{0.1v + 4}{1 - \exp\left(-0.1v - 4\right)},\tag{9}$$

$$\beta_m(v) = 4 \exp\left(\frac{-v - 65}{18}\right),\tag{10}$$

$$\alpha_h(v) = 0.07 \exp\left(\frac{-v - 65}{20}\right),\tag{11}$$

$$\beta_h(v) = \frac{1}{1 + \exp\left(-0.1v - 3.5\right)},\tag{12}$$

where v = V/[mV] is the membrane voltage V of Eq. (1) in millivolts (mV) divided by 1mV to render the variable dimensionless. We consider $C = 1\mu\text{F/cm}^2$, $g_{\text{K}} = 36\text{mS/cm}^2$, $g_{\text{Na}} = 120\text{mS/cm}^2$, $g_{\text{L}} = 0.3\text{mS/cm}^2$, $E_{\text{K}} = -77\text{mV}$, $E_{\text{Na}} = 50\text{mV}$, $E_{\text{L}} = -54.4\text{mV}$ and $V_r = 20\text{mV}$.

Figure 1(a) shows the spike (or natural) frequency ν (Hz) of a single HH neuron as a function of current *I*. The spikes were numerically computed when the voltage *V* crosses the threshold of 0mV, increasing from negative to positive values. In the simulations of 100 coupled HH neurons in Sec. III, we consider I_i randomly distributed in the interval $[10, 30]\mu A/cm^2$, leading to spike frequencies ν_i in the interval [70, 100]Hz. This interval avoids the regime observed for *I* in $[6, 10]\mu A/cm^2$, as shown in Fig. 1(b), where it is clear that for some *I* values, the neuron does not spike as its natural frequency ν is zero. This choice of interval also allows for the monotonic increase of the natural frequencies ν_i without reaching currents *I* bigger than $60\mu A/cm^2$ that correspond to a non-spike regime.

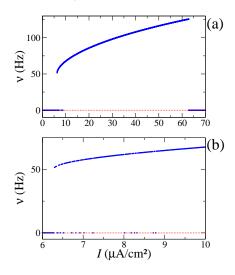


FIG. 1. Natural frequency ν as a function of the current I for a single HH neuron. Panel (b) is a zoom-in of panel (a) for currents I in [6, 10] μ A/cm².

In Eq. (5), f_i is the strength of the effective synaptic output current from neuron i to neuron j and τ_s (ms) the synaptic time constant, fixed at $\tau_s = 2.728$ ms. When neuron ispikes, f_i is updated $(f_i \to 1)$ before it starts to decay exponentially. Equation (6) models STP^{33,50} with τ_D (ms) being the recovery time constant, related to

Equation (6) models STP^{33,50} with τ_D (ms) being the recovery time constant, related to biological mechanisms such as the depletion of release-ready neurotransmitter vesicles at the presynaptic terminal^{44,51}. We assume that every time neuron *i* spikes, the update rule

 $D_i \rightarrow (D_i - d)$ is applied. Biologically, D_i could represent the vesicles that can be used to transmit a signal from the presynaptic to the postsynaptic neuron. Therefore, the update rule means that the amount D_i of available vesicles is decreased by d = 0.1 in every spike of neuron *i*, and then it recovers according to Eq. (6). In that framework, D_i lies in [0, 1] as if it happens D_i to be negative it means that the neuron used all the stored vesicles and when this happens, D_i is reset to 0. On the other hand, when D_i is equal to 1, it means that all neurotransmitter vesicles are restored.

Figure 2 shows the effect of STP on a pair of neurons coupled with a unidirectional connection from neuron 1 to neuron 2. This is implemented by fixing $\varepsilon_{12} = 0$ (i.e. the connectivity strength from neuron 2 to neuron 1 is 0) and $\varepsilon_{21} = 0.1$ for the connectivity strength from neuron 1 to 2. In this study, we have set the STP recovery time τ_D at 50ms and I_2 at 0, so that neuron 2 spikes only when it receives a strong enough input $I_{2,input} =$ $(V_r - V_2)\varepsilon_{21}f_1D_1$ from neuron 1. When the spiking frequency ν_1 in Fig. 2(a) changes from 70Hz to 100Hz, ν_2 in Fig. 2(b) exhibits a delayed alteration in its dynamic behaviour as its amplitude plummets at about 300ms. We appreciate further this phenomenon through the temporal evolution of D_1 in Fig. 2(c) and the input current $I_{2,\text{input}}$ received by neuron 2 in Fig. 2(d). Comparing the 70Hz-regime with the 100Hz-regime in Figs. 2(c) and (d), one can see that D_1 decreases with the increase of the spike-frequency of neuron 1 to 100Hz, and consequently $I_{2,\text{input}}$ becomes less intense as it is not strong enough to cause spikes in the activity of neuron 2. When neuron 1 returns to 70Hz spike frequency, there is more time for D_1 to recover, thereby to increase the intensity of $I_{2,input}$ which triggers again spikes in neuron 2. Generalizing this, one might say that STP makes neurons more sensitive to spike frequency changes.

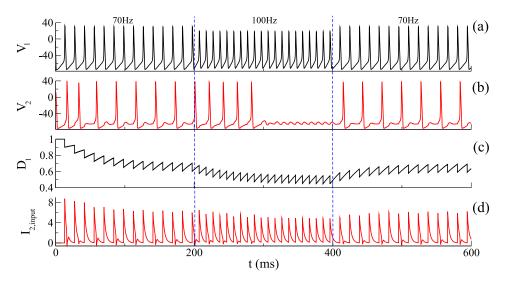


FIG. 2. The effect of STP on a pair of unidirectionally connected HH neurons, where neuron 1 is connected to neuron 2 but not vice versa: Temporal evolution of (a) ν_1 , (b) ν_2 , (c) D_1 and (d) $I_{2,input} = (V_r - \nu_2)\varepsilon_{21}f_1D_1$ with $\varepsilon_{12} = 0$, $\varepsilon_{21} = 0.1$, $I_2 = 0$ and $\tau_D = 50$ ms. Note that, for $\nu_1 = 70$ Hz and $\nu_1 = 100$ Hz, we used $I_1 = 10.97\mu$ A/cm² and $I_1 = 31.8\mu$ A/cm², respectively.

Moving now to the other form of plasticity introduced to the model, STDP gives rise to changes in the synaptic strength by means of the update function^{39,48}

$$\varepsilon_{ij} \to \varepsilon_{ij} + 10^{-3} \Delta \varepsilon_{ij},$$
 (13)

where

$$\Delta \varepsilon_{ij} = \begin{cases} \varepsilon_+ = A_1 e^{(-\Delta t_{ij}/\tau_1)}, & \text{if } \Delta t_{ij} > 0\\ \varepsilon_- = -A_2 e^{(\Delta t_{ij}/\tau_2)}, & \text{if } \Delta t_{ij} < 0\\ 0, & \text{if } \Delta t_{ij} = 0 \end{cases}$$
(14)

Here, $\Delta t_{ij} = t_i - t_j$ is the difference between the spike times of the postsynaptic (t_i) and presynaptic (t_j) neurons *i* and *j*, respectively. Figure 3 shows the plot of the plasticity function $\Delta \varepsilon_{ij}$ calculated from Eq. (14) for $A_1 = 1$, $A_2 = 0.5$, $\tau_1 = 1.8$ ms, $\tau_2 = 6$ ms, and Δt_{ij} varying from -20ms to 20ms. This update rule is applied everytime the postsynaptic neuron *i* spikes.

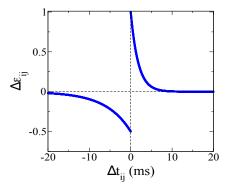


FIG. 3. Plot of $\Delta \varepsilon_{ij}$ that models STDP as a function of the difference between the spike times of the postsynaptic neuron *i* and presynaptic neuron *j*. Note the discontinuity at $\Delta t_{ij} = 0$ ms.

To understand better the effects of the recovery time τ_D , we consider in the next the case of two, unidirectionally connected neurons where neuron 1 (presynaptic) is connected to neuron 2 (postsynaptic), but not vice versa. We evaluate how the average input current $\bar{I}_{2,\text{input}}$ and amount of information exchanged change with the increase of the recovery time τ_D by comparing cases where the neurons have similar and dissimilar spike frequencies.

To quantify the exchange of information between the 2 neurons, we compute the Mutual Information $(MI)^{54-56}$. MI can be understood as the amount of dependence (or uncertainty) between two random variables X and Y, and is given by

$$\mathrm{MI}_{XY}(n) = \sum_{i}^{n} \sum_{j}^{n} P_{XY}(i,j) \log\left(\frac{P_{XY}(i,j)}{P_{X}(i)P_{Y}(j)}\right)$$
(15)

where the probability of a random event i(j) to occur in X(Y) is given by $P_X(i)(P_Y(j))$. The joint probability $P_{X,Y}(i,j)$ gives the probability of i to occur in X and j in Y, simultaneously. In this context, the number of random events in X and Y is denoted by n.

In the next, we consider the cases where the system has only STP (see Fig. 4(a), (c)) and, STP and STDP (see Fig. 4(b), (d)). In both cases, we consider $\varepsilon_{21} = 0.3$ and $\varepsilon_{12} = 0$, but for the case of STD and STDP (Fig. 4(b), (d)), STDP acts only on ε_{21} (unidirectional connection) and ensures that the spike frequency of presynaptic neuron 1 remains unchanged. The spike frequency of neuron 2 was fixed at $\nu_2 = 70$ Hz ($I_2 = 10.97\mu$ A/cm²) and, for the blue curves in Fig. 4, neuron 1 has $\nu_1 = 100$ Hz ($I_1 = 31.8\mu$ A/cm²) and for the orange curves, $\nu_1 = 72$ Hz ($I_1 = 11.88\mu$ A/cm²). We use as random variables to calculate MI, the time series of the voltage of both neurons, namely $X = V_1$ and $Y = V_2$. We let the system evolve for 200×10^3 ms for each τ_D and used the last 100×10^3 ms to calculate MI. For the calculation of $\bar{I}_{2,input}$, we consider the last 10×10^3 ms of the simulations.

Observing Fig. 4(a) (STP), one can see that for recovery times $\tau_D < 75$ ms, the average current $\bar{I}_{2,\text{input}}$ for $\nu_1 = 100$ Hz (blue curve) is bigger than $\bar{I}_{2,\text{input}}$ for $\nu_1 = 72$ Hz (orange curve) (see also the inset in Fig. 4(a)). Interestingly, this change in the region 75ms< $\tau_D < 125$ ms where $\bar{I}_{2,\text{input}}$ for $\nu_1 = 72$ Hz (orange curve) is bigger than $\bar{I}_{2,\text{input}}$ for $\nu_1 = 100$ Hz (blue curve). For $\tau_D > 125$ ms, both $\bar{I}_{2,\text{input}}$ curves assume similar values and settle asymptotically to zero with further increasing in τ_D .

In Fig. 4(b) (STP and STDP), we see that the action of STP and STDP causes $\bar{I}_{2,\text{input}}$ for $\nu_1 = 100$ Hz to drop to 0 at $\tau_D \approx 100$ ms whereas $\bar{I}_{2,\text{input}}$ for $\nu_1 = 72$ Hz (orange curve) stays

positive until $\tau_D \approx 480$ ms. Panels (c) and (d) show clearly the impact of STP and STDP to MI. In particular, MI fluctuates around 3 for $\tau_D < 480$ ms in both panels. However, both panels show that as τ_D increases from 0ms, MI for $\nu_1 = 100$ Hz (blue curve) approaches zero at $\tau_D \approx 100$ ms, while for $\nu_1 = 72$ Hz, it remains close to 3 until $\tau_D \approx 480$ ms. These results show that STP make the influence of slower (or close frequency) neurons to become greater than that of faster (or very different frequency) neurons as τ_D increases. In conjunction with STDP, this effect occurs due to the decrease of $\bar{I}_{2,input}$ that makes the firing times uncorrelated (less synchronized), which then causes the coupling to disappear, i.e. $\varepsilon_{21} \rightarrow 0$. This decoupling process tends to occur for smaller τ_D values as the difference between ν_1 and ν_2 is amplified. We address this further later when we discuss the results in Fig. 5.

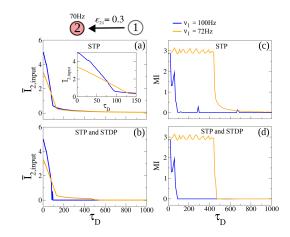


FIG. 4. The average input current $\bar{I}_{2,\text{input}}$ neuron 2 receives from neuron 1 and mutual information, MI, as a function of the recovery time τ_D for the following cases: (a), (c) only STP and (b), (d) STP and STDP. Neuron 2 has a fixed spike frequency $\nu_2 = 70$ Hz ($I_2 = 10.97\mu$ A/cm²). The blue curve represents the case where $\nu_1 = 100$ Hz ($I_1 = 31.8\mu$ A/cm²) and the orange curve, the case where $\nu_1 = 72$ Hz ($I_1 = 11.88\mu$ A/cm²). We use $\varepsilon_{21} = 0.3$ and $\varepsilon_{12} = 0$, implying that only ε_{21} changes when STDP is considered. The simulations run for 200 × 10³ms, $\bar{I}_{2,\text{input}}$ is calculated over the last 10 × 10³ms and, the time-series for ν_1 and ν_2 used for the calculation of MI are recorded from the last 100 × 10³ms of the simulations.

III. EFFECTS OF STP ON HH NEURAL NETWORKS WITH STDP

Neural networks with STDP and random synaptic input were studied by Popovych et al.³⁹. The authors reported that the mean synaptic coupling depends on the noise intensity. Recently, the authors in²⁰ showed that STDP induces non-trivial topology in neural networks. Here, we extend this work and build a neural network of N = 100 HH neurons to study the combined effect of STP and STDP on the structure of the network and in particular, on its connectivity.

We start by analysing neural connectivity under the effect of STP and STDP. To this end, we consider a pair of HH neurons with STP and STDP, with the coupling strengths ε_{12} and ε_{21} varying in [0, 0.3], and the connectivity threshold set at 0.01. This threshold is chosen because the coupling weights that should vanish in time actually oscillate around 0 assuming very small values as STDP is always present and thus, affecting them. As a consequence, the time averages of the coupling weights are not 0 but very close to 0. Moreover, for coupling weights smaller than this threshold, we notice that neurons influence each other only slightly.

Initially, the pair of neurons is uncoupled (i.e. $\varepsilon_{12} = \varepsilon_{21} = 0$) or bidirectionally coupled with $\varepsilon_{12} = \varepsilon_{21} = 0.3$. Figure 5 shows the direction of connectivity, after a transient time, for different natural frequencies ν_1 and ν_2 in [70, 100]Hz. The direction is coloured

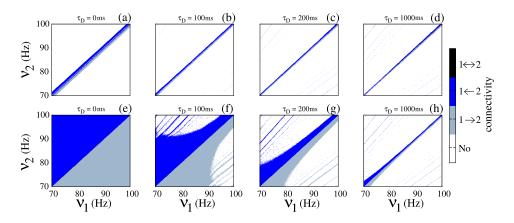


FIG. 5. Parameter spaces $\nu_1 \times \nu_2$ for a pair of initially uncoupled neurons for recovery times (a) $\tau_D = 0$ ms, (b) $\tau_D = 100$ ms, (c) $\tau_D = 200$ ms and (d) $\tau_D = 100$ ms, and for a pair of initially coupled neurons for (e) $\tau_D = 0$ ms, (f) $\tau_D = 100$ ms, (g) $\tau_D = 200$ ms and (h) $\tau_D = 100$ 0ms. Note that ν_1 and ν_2 vary in [70, 100]Hz and that the colour bar shows the direction of synaptic connectivity, where white accounts for the uncoupled case (denoted "No"), grey for the case where neuron 1 is connected to neuron 2 (i.e. $1 \rightarrow 2$), blue for the case where neuron 2 is connected to neuron 1 (i.e. $1 \leftarrow 2$) and black for the bidirectional connection (i.e. $1 \leftrightarrow 2$).

according to the connection from the output of one neuron to the input to the other neuron (directions of arrows in the colour bar): white accounts for the uncoupled case denoted in the colour bar as "No" ($\bar{\epsilon}_{ij} < 0.01$), grey for the case where neuron 1 is connected to neuron 2 (i.e. $1 \rightarrow 2$), blue for the case where neuron 2 is connected to neuron 1 (i.e. $1 \leftarrow 2$), and black for the bidirectional connection (i.e. $1 \leftrightarrow 2$). For initially uncoupled neurons, the connections remain only for neurons with very similar frequencies, from those with faster spike frequencies to those with slower as depicted in Fig. 5(a) ($\tau_D = 0$ ms), 5(b) ($\tau_D = 100$ ms), 5(c) ($\tau_D = 200$ ms), and 5(d) ($\tau_D = 1000$ ms). We note the absence of bidirectional connections and that the increase of the recovery time τ_D gives rise to a narrower region of directed connections. With regard to neurons starting with bidirectional connections, we observe that for $\tau_D = 0$ ms, neurons still remain unidirectionally connected (for all natural frequencies) from the high to the low frequency neurons (see Fig. 5(e)). Increasing τ_D to 100ms (see Fig. 5(f)), 200ms (Fig. 5(g)) and 1000ms (Fig. 5(h)), we observe a decrease in the area that represents connectivity, implying that the region of high spike frequencies is more affected by the influence of STP. Therefore, STDP makes the connections increase from faster spiking to slower spiking neurons and STP decreases the influence of hight frequency neurons, allowing connections only for those neurons with similar spike frequencies.

The reason for the initially, bidirectionally coupled system to have larger areas of connectivity is related to the difference in frequency $\Delta \nu_{ij}^c$ of neurons when coupled (or the synchronization level). We note that this difference is not equal to the difference between their natural frequencies $\nu_i - \nu_j$, i.e. $\Delta \nu_{ij}^c \neq \nu_i - \nu_j$. In Fig. 6, we present the calculation of the average $\Delta \nu_{ij}^c$ (colour bar) for the first t = 4000ms for the system with STP only. We consider N = 2, $\varepsilon_{12} = \varepsilon_{21} = 0.3$ and vary ν_1 and ν_2 in [70, 100]Hz. Figure 6(a) shows the result for $\tau_D = 0$ ms where, despite the natural frequencies of the neurons, $\Delta \nu_{ij}^c$ is approximately equal to 0 (black region). When comparing these results with the results in Fig. 5(e), once can see that this corresponds to a connected region. For increasingly bigger recovery times τ_D (see Fig. 6(b) for $\tau_D = 100$ ms, 6(c) for $\tau_D = 200$ ms and 6(d) for $\tau_D = 1000$ ms), one can observe a decrease in the size of the black area, similar to the size of the area seen in Fig. 5(f), (g) and (h). Thus, one can infer that what defines a connected configuration in the case of STDP is the difference in spike frequency among neurons or how synchronized they become. The non-homogeneous distribution of connected areas in relation to the main diagonal in Figs. 5 occurs due to the non-linear variation of neural

frequencies in relation to the received external current. In Fig. 1, we see that for low currents (i.e. $I \in [10, 25]$), a small increase in I can cause a bigger variation in the frequencies when compared to the interval where I > 30. Therefore, low frequency neurons are more sensitive to changes in their external currents, which facilitates their synchronization with neurons with similar frequencies. The more synchronized they become, the more they remain connected. Coming back to the effect of STP, we find that it decreases the influence of the faster neurons on the slower ones, and consequently leads to the increase of their frequency differences and to the suppress of their synchronization. That then leads to the depression of the connectivity via STDP.

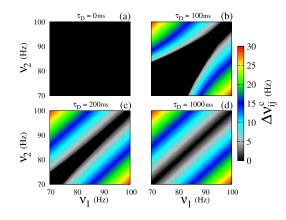


FIG. 6. Parameter space $\nu_1 \times \nu_2$ for increasing values of τ_D : (a) for $\tau_D = 0$ ms, (b) for $\tau_D = 100$ ms, (c) for $\tau_D = 200$ ms and (d) for $\tau_D = 1000$ ms. Note that ν_1 and ν_2 vary in [70, 100]Hz and the colour bar represents $\Delta \nu_{ij}^c$ values.

We performed a similar study for a HH neural network and analyzed how the connections evolve with the combined application of STP and STDP. In particular, we start with an all-to-all (globally connected) network of N = 100 excitatory, HH neurons. The coupling weights vary in [0, 0.04] and I_i is randomly distributed so that neural frequencies are in the range [70, 100] Hz. We choose three initial coupling weight averages $\bar{\varepsilon}_{\text{initial}}$. The coupling matrices for the connectivity strengths ε_{ij} are shown in Fig. 7 where $\overline{\varepsilon}_{initial} = 0$ in 7(a), $\bar{\varepsilon}_{\text{initial}} = 0.01 \text{ in } 7(\text{b}) \text{ and } \bar{\varepsilon}_{\text{initial}} = 0.04 \text{ in } 7(\text{c}).$ For all initial conditions studied, we kept the standard deviation of $\bar{\varepsilon}_{initial}$ fixed at 0.002. In the coupling matrix (ε_{ij}), the presynaptic neurons j and postsynaptic neurons i are sorted and plotted in ascending frequency-order (i.e. from the smallest to the largest spike frequency). The final coupling matrices ε_{ij} for $\tau_D = 0$ ms are shown in Fig. 7(d), (e) and (f) for $\bar{\varepsilon}_{initial} = 0$, $\bar{\varepsilon}_{initial} = 0.01$ and $\bar{\varepsilon}_{initial} = 0.04$, respectively. In all cases, the final coupling matrices have, predominantly, connections from faster to slower spiking neurons. This behaviour was also reported in Borges et $al.^{20}$ for a neural network with STDP. Figure 7(g), (h) and (i) present our results for $\bar{\varepsilon}_{\text{initial}} = 0$, $\bar{\varepsilon}_{\text{initial}} = 0.01$, and $\bar{\varepsilon}_{\text{initial}} = 0.04$, respectively, where $\tau_D = 100$ ms. Due to the effect of STP on the dynamics of the neurons in the network, we observe the formation of different modules of directly connected neurons. Again, the effect of STP leads to a decrease on the influence of the fastest neurons to the slowest ones, allowing for the formation of connections among those neurons with similar spike frequencies. The size of these modules increases according to the intensity of the initial coupling and their different sizes can be explained by the analysis made based on the results in Figs. 5 and 6, where the coupling shortened the frequency differences, leading to the formation of connections especially among neurons with close frequencies. It also explains why the bigger-size modules are composed of the smallest-frequency spiking neurons. Increasing τ_D to 1000ms, leads to the disappearance of big modules and to the decrease of the number of coupled neurons, as shown in Fig. 7(j), (k) and (l). Therefore, by varying the STP recovery time τ_D , one can control the formation of modules in neural networks with STP and STDP.

To study further the observed modules, we consider connections with $\varepsilon_{ij} > 0.002$ (i.e.

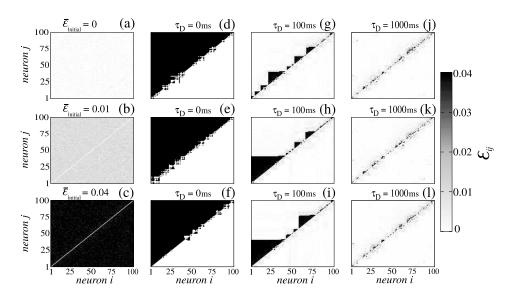


FIG. 7. The combined effect of STP and STDP on connectivity in networks of N = 100 excitatory coupled HH neurons and emergence of modular neural networks. The initial couplings used are: In (a) $\bar{\varepsilon}_{\text{initial}} = 0$, in (b) $\bar{\varepsilon}_{\text{initial}} = 0.01$ and in (c) $\bar{\varepsilon}_{\text{initial}} = 0.04$. We consider $\tau_D = 0$ ms for (d) $\bar{\varepsilon}_{\text{initial}} = 0$, (e) $\bar{\varepsilon}_{\text{initial}} = 0.01$ and (f) $\bar{\varepsilon}_{\text{initial}} = 0.04$. $\tau_D = 100$ ms for (g) $\bar{\varepsilon}_{\text{initial}} = 0$, (h) $\bar{\varepsilon}_{\text{initial}} = 0.01$ and (i) $\bar{\varepsilon}_{\text{initial}} = 0.04$ and $\tau_D = 1000$ ms for (j) $\bar{\varepsilon}_{\text{initial}} = 0$, (k) $\bar{\varepsilon}_{\text{initial}} = 0.01$ and (l) $\bar{\varepsilon}_{\text{initial}} = 0.04$. Note that the synaptic weights ε_{ij} (with i, j = 1, ..., 100) of the coupling matrices are encoded in grey scale in the colour bar.

with connectivity strength bigger than 5% of the maximal coupling strength). Thus, weaker connections are not considered in the resulting network analysis. This procedure avoids measurement errors that might be caused by connections whose weights fluctuate over time closely to zero.

To evaluate how modular structures evolve over time, we compute the modularity Q by using the Louvain method⁵⁸. Q is measured across network partitions in densely connected communities. In particular, the modularity assumes values in the range [-1,1] comparing the density of connections within communities with the density among communities. The best network partition in modules is one that maximizes modularity. Q is defined as⁵⁹

$$Q = \frac{1}{W} \sum_{i}^{N} \sum_{j}^{N} \left(\varepsilon_{i,j} - \frac{\omega_i \omega_j}{W} \right) \delta(c_i, c_j), \tag{16}$$

where $\omega_i = \sum_j^N \varepsilon_{ij}$ represents the sum of the connection weights received by node *i* and $W = \sum_i^N \sum_j^N \varepsilon_{ij}$ is the sum of all weights in the coupling matrix. The term c_i represents the community that neuron *i* has been allocated to and $\delta(c_i, c_j)$ is given by

$$\delta(c_i, c_j) = \begin{cases} 1, \text{ if } c_i = c_j, \\ 0, \text{ otherwise.} \end{cases}$$
(17)

The Louvain⁵⁸ method is defined in two steps. At first, each node in the network is considered as a community in itself, and thus initially, there will be as many communities as nodes in the network. At this stage, each node i is reassigned to the community of each of its neighbors j, then i will be permanently fixed in the community that promotes the largest gain in modularity Q (positive gain). This process is applied repeatedly to all network nodes until there are no more gains in Q. The second step amounts to taking the defined communities at the end of the first step and consider them as the nodes of a new network. The weight of the connections between these new nodes is given by the sum of the weights of the connections between the nodes present in the communities (defined in the first step). Once the new network is computed, the first step can be applied again to its nodes. This process occurs repeatedly until no further changes to Q occur and a maximum value is thus obtained.

Figure 8 shows the time evolution of the modularity Q for the three recovery times, $\tau_D = 0$ ms, $\tau_D = 100$ ms and $\tau_D = 100$ ms. The colours represent the initial connectivity strengths in the coupling matrix: the black curve is for $\bar{\varepsilon}_{initial} = 0$, the red curve for $\bar{\varepsilon}_{\text{initial}} = 0.01$ and the green curve for $\bar{\varepsilon}_{\text{initial}} = 0.04$. The coloured symbols represent the average Q values calculated over 20 random networks obtained by rewiring all connections of the original networks. We did this to compare our results with those obtained for random networks with the same connections and number of nodes (which we call random variants). Unless stated otherwise, the symbols in Fig. 8 represent the measurements taken on these random variants. In Fig. 8(a), we observe that for $\tau_D = 0$ ms (instantaneous recovery), Q is very low for all $\bar{\varepsilon}_{\text{initial}}$ connectivity strengths, remaining constant after a small transient. These results confirm those in Fig. 7(d), (e) and (f). In Fig. 8(b) for $\tau_D = 100$ ms, the modularity Q takes its largest value for $\bar{\varepsilon}_{\text{initial}} = 0$, confirming what we observed by comparing the number of modules in Fig. 7(g), (h) and (i). It is also apparent that Q stabilizes for simulation times t greater than 80000ms as there are no changes in the modular structures occurring in the networks anymore. Interestingly enough, Fig. 8(c) shows that for $\tau_D = 1000$ ms, the modularity converges to the same value (i.e. $Q \approx 0.6$) for all $\bar{\varepsilon}_{initial}$ coupling strengths, again in accordance with the results in Fig. 7(j), (k), (l), which show networks with similar configurations regardless of the initial coupling. Finally, in all cases considered, we find that the modularity of the networks is bigger than their random variants.

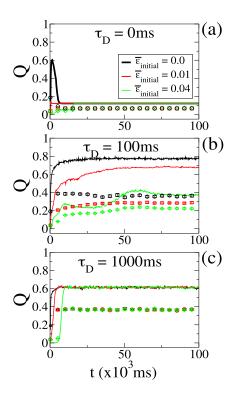


FIG. 8. Time evolution of the modularity Q for (a) $\tau_D = 0$ ms, (b) $\tau_D = 100$ ms and (c) $\tau_D = 100$ oms. Q was calculated over the time evolution of the coupling matrix and the coloured curves represent the network average initial coupling: the black curve is for $\bar{\varepsilon}_{\text{initial}} = 0$, the red for $\bar{\varepsilon}_{\text{initial}} = 0.01$ and the green for $\bar{\varepsilon}_{\text{initial}} = 0.04$. The colored symbols represent the average Q for random networks computed by rewiring the connections in the corresponding original networks. For example, the black symbols represent the average Q calculated over 20 random networks generated from rewiring randomly the connections in the original networks (black curve).

Next, we compute other quantities that characterize the structure of the networks considered previously, such as the mean path-length, clustering coefficient and assortativity.

In particular, a path is defined as the route that passes through network connections connecting two nodes i and j. The path with the shortest number of connections is called the shortest path l_{ij}^{60} . The networks we consider here are directed, so l_{ij} is not necessarily equal to l_{ji} . In general, the shortest average path length is given by

$$\bar{L} = \frac{1}{N(N-1)} \sum_{i}^{N} \sum_{j \neq i}^{N} l_{ij},$$
(18)

where N is the number of nodes in the network. In our approach, we calculated \bar{L} via a breadth-first search approach⁶⁰ and do not consider the weights of the connections. If there is no possible directed path l_{ij} between nodes *i* and *j*, then it is not considered in the calculations in Eq. (18).

The second computed quantity is the clustering coefficient CC_i , which measures the degree to which the neighbors of a node *i* are connected to each other and varies in [0, 1]. It is computed by considering the number of triangular motifs made by node *i* and its neighbors compared to all possible triangular motifs of that node⁶⁰. For directed networks, given 3 connected nodes *i*, *j* and *h*, there are 8 distinct triangular motifs, shown in Fig. 9. These motifs are further organised into 4 groups when considering node *i* as the reference node: Figure 9(a) shows a "cycle" motif, Fig. 9(b) a "middleman" motif, Fig. 9(c) an "in" motif and Fig. 9(d) an "out" motif. For each motif, we calculate $CC_i^{(cyc,mid,in,out)}$ relative to node *i*, as shown in⁶¹. For directed and weighted networks, there are 4 types of clustering coefficients

$$CC_{i}^{Cyc} = \frac{\frac{1}{2} \sum_{j}^{N} \sum_{h}^{N} \left[\varepsilon_{ij}^{1/3} \varepsilon_{jh}^{1/3} \varepsilon_{hi}^{1/3} + \varepsilon_{ih}^{1/3} \varepsilon_{hj}^{1/3} \varepsilon_{ji}^{1/3} \right]}{d_{i}^{in} d_{i}^{out} - d_{i}^{\leftrightarrow}},$$

$$CC_{i}^{mid} = \frac{\frac{1}{2} \sum_{j}^{N} \sum_{h}^{N} \left[\varepsilon_{ih}^{1/3} \varepsilon_{jh}^{1/3} \varepsilon_{ji}^{1/3} + \varepsilon_{ij}^{1/3} \varepsilon_{hi}^{1/3} \varepsilon_{hj}^{1/3} \right]}{d_{i}^{in} d_{i}^{out} - d_{i}^{\leftrightarrow}},$$

$$CC_{i}^{in} = \frac{\frac{1}{2} \sum_{j}^{N} \sum_{h}^{N} \left[\varepsilon_{ij}^{1/3} \varepsilon_{ih}^{1/3} \varepsilon_{jh}^{1/3} + \varepsilon_{ij}^{1/3} \varepsilon_{ih}^{1/3} \varepsilon_{hj}^{1/3} \right]}{d_{i}^{in} (d_{i}^{in} - 1)},$$

$$CC_{i}^{out} = \frac{\frac{1}{2} \sum_{j}^{N} \sum_{h}^{N} \left[\varepsilon_{hi}^{1/3} \varepsilon_{ji}^{1/3} \varepsilon_{jh}^{1/3} + \varepsilon_{hi}^{1/3} \varepsilon_{ji}^{1/3} \varepsilon_{hj}^{1/3} \right]}{d_{i}^{out} (d_{i}^{out} - 1)},$$
(19)

where $d_i^{\text{in}} = \sum_j^N a_{ij}$ is the in-degree of node *i* and $d_i^{\text{out}} = \sum_j^N a_{ji}$ its out-degree. The term $d_i^{\leftrightarrow} = \sum_j^N a_{ij}a_{ji}$ represents the number of bilateral connections between node *i* and its neighbors. For the calculation of d_i^{in} , d_i^{out} and d_i^{\leftrightarrow} , the coupling weights are not considered, only the number of connections, that is $a_{ij} = 1$ if $\varepsilon_{ij} > 0.002$, otherwise $a_{ij} = 0$.

Consequently, the clustering coefficient of the network, CC^* , is calculated by averaging CC_i^* over all N nodes in the network

$$CC^* = \frac{1}{N} \sum_{i}^{N} CC_i^*, \qquad (20)$$

where * stands for either of the cyc, mid, in or out motifs.

The last quantity in our study is assortativity, which is the correlation coefficient (i.e. the Pearson correlation⁶²) between the degrees of nodes on two opposite ends of a connection in a network, for all connections in the network. This correlation varies in [-1, 1] and, is

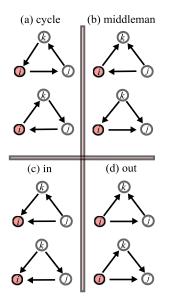


FIG. 9. The 8 distinct triangular motifs in directed networks. The motifs are considered with respect to node i depicted in red and are split into the cycle, middleman, in and out motifs.

positive in assortative networks and negative in disassortative networks. Since our networks are directed and weighted, we use the four directed assortativity measures defined in⁶³.

In particular, let a and b be indices representing the type of degree (i.e. in- and outdegree) and j_i^a and k_i^b the a- and b-degree from source node j to target node k of an edge e. Figure 10 shows an illustrative representation of edge e and their respective source nodes jand target nodes k. Foster et al.⁶³ defined the assortativity measure r(a, b) as

$$r(a,b) = \frac{N_e^{-1} \sum_{e}^{N_e} \left[(j_e^a - \bar{j}^a) (k_e^b - \bar{k}^b) \right]}{\sigma^a \sigma^b},$$
(21)

where N_e is the number of edges in the network, $\overline{j}^a = N_e^{-1} \sum_e^{N_e} j_e^a$ the average in- and out-degrees of the source node and, $\sigma^a = \sqrt{N_e^{-1} \sum_e^{N_e^{-1}} (j_e^a - \overline{j}^a)^2}$ its deviation, calculated for all edges. The equations for \overline{k}^b and σ^b are similarly defined. In our approach, instead of in- and out-degrees, we use the in- and out-strength from the source and target nodes, given by the sum of their in- and out-coupling weights. A positive assortativity coefficient indicates that nodes tend to connect to other nodes with the same or similar strength.

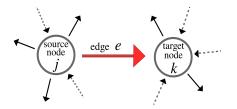


FIG. 10. Representative illustration of an edge e and, the source j and target k nodes considered in Eq. (21).

Figure 11 shows the results of the computations of the quantities discussed previously as a function of τ_D in [0, 1000]ms. In all computations, we consider the final coupling matrix obtained after t = 200000ms of simulated data. Figure 11(a) shows the modularity Q for 3 initial couplings: $\bar{\varepsilon}_{initial} = 0$ in black, $\bar{\varepsilon}_{initial} = 0.01$ in red and, $\bar{\varepsilon}_{initial} = 0.04$ in green and their random variants given by different symbols with their respective colors. We note that for $\tau_D < 30$ ms (approximately) and for all 3 $\bar{\varepsilon}_{initial}$ values, Q is low. For 30ms $< \tau_D < 120$ ms (approximately), Q depends on $\bar{\varepsilon}_{initial}$. In particular, the lower the initial average coupling $\bar{\varepsilon}_{initial}$, the more modular the final configuration of the network. For $\tau_D > 120$ ms (approximately), there is no dependence on $\bar{\varepsilon}_{initial}$ and the resulting networks exhibit similar modularity values. In all cases of $\bar{\varepsilon}_{initial}$, Q is higher than that calculated for their random variants, which shows that the resulting networks do not have characteristics of random networks.

Figure 11(b) shows the results for the average path length \bar{L} for the same three $\bar{\varepsilon}_{\text{initial}}$ values. For $\tau_D < 30$ ms, the average path $\bar{L} = 1$, in accordance with Fig. 7(d), (e) and (f), where we observe a network where all nodes are connected to all other nodes with unidirectional connections. For 30ms $< \tau_D < 120$ ms, the average path length \bar{L} is bigger for weaker initial couplings. This is because weak initial couplings result in the formation of a larger number of modular structures (as seen in Fig. 11(a)), which results in a greater path for one node to access other nodes in different modules, possibly crossing through other modules. For $\tau_D > 120$ ms, we see that $\bar{L} \approx 2.5$ for all 3 $\bar{\varepsilon}_{\text{initial}}$ values. In contrast, the average path lengths \bar{L} for the random variants are smaller than those for the original networks.

Figure 11(c), (d) and (e) present the 4 types of clustering coefficients for $\bar{\varepsilon}_{initial} = 0$, $\bar{\varepsilon}_{\text{initial}} = 0.01$ and $\bar{\varepsilon}_{\text{initial}} = 0.04$. The black curve represents CC^{cyc} , the red CC^{mid} , the green CC^{in} and the blue CC^{out} . In all cases, the results are similar. It is worth it to note that for $\tau_D < 120$ ms, CC^{mid} has the highest value and CC^{cyc} the lowest. These results build on what we have already observed: the action of STDP promotes connections from faster to slower spiking neurons and do not permit cyclic connections (see Fig. 9(a)). As τ_D increases, CC^{mid} converges to CCⁱⁿ, CC^{out} ≈ 0.15 and CC^{cyc} increases to ≈ 0.09 . For the random variants of the networks (coloured symbols), we see that for $\tau_D < 120$ ms, all clustering coefficients are similar and larger than those for the original networks. This is because for such τ_D values, the networks from which they were generated (Fig. 7(d), (e), (f)) are densely connected and the high number of connections allows for the formation of triangular motifs without any of the 4 types occuring preferentially. For $\tau_D > 120$, all random networks have their clustering coefficients fluctuate near zero. In this case, the generating networks are more sparse (as in Fig. 7(g), (h), (i)), which makes their random variants have low probability in forming triangular motifs. We note here that these clustering coefficients are not high enough to claim with certainty that the networks have a small-world topology.

The right column of plots in Fig. 11 shows the 4 assortativity measures, r, computed for the considered networks with $\bar{\varepsilon}_{\text{initial}} = 0$, $\bar{\varepsilon}_{\text{initial}} = 0.01$ and $\bar{\varepsilon}_{\text{initial}} = 0.04$. The black curves represent the out-strength/in-strength correlation r(out, in) between the source and target nodes, the red curve the in-strength/out-strength correlation r(in, out), the green curve the out-strength/out-strength correlation r(out, out) and the blue curve, the in-strength/instrength correlation r(in, in). In Fig. 11(f), (g), (h) and for the three $\bar{\varepsilon}_{initial}$ values, we find that for $\tau_D < 100$ ms, the networks are disassortative as r(out, in) and r(in, out) are negative and, at the same time, neither assortative nor disassortative as r(out, out) and r(in, in) are approximately equal to 0. As τ_D increases, all correlations grow being mostly positive with r(out, in) being the largest. This shows a greater correlation in the network of high out-strength nodes to connect with nodes with high in-strength and corroborates the results in Fig. 7 where connections occur preferably from the faster neurons (which have a high out-strength) to the slower ones (with high in-strength). For the random variants of these networks, all correlations are close to zero for the entire τ_D range, a result completely different to those for the original networks. We thus conclude that the original networks are far from being purely random networks, exhibiting a type of preferential attachment in their connectivities.

In the analysis of the structural properties of the networks, the topology may vary greatly depending on τ_D . In all cases studied, the networks were different from random or small-world networks, since the average path lengths are bigger than in random networks and

their clustering coefficients are small than in random networks. This corroborates further the results obtained previously here that there is a complex configuration where modules are formed following a kind of preferential attachment process.

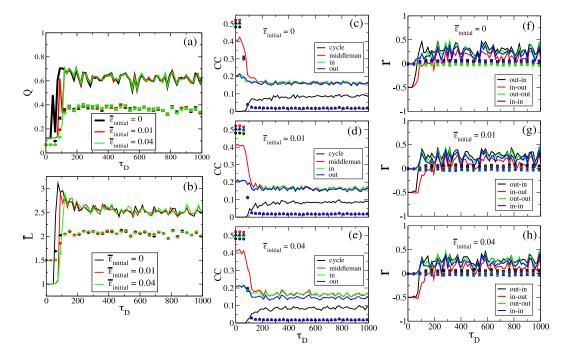


FIG. 11. Structural properties of the final coupling matrix configuration (coloured curves) and their respective random variants (averaged over 20 random networks, coloured symbols) as a function of the recovery time τ_D . Panels (a) and (b) show the modularity Q and average shortest path length \bar{L} for three different average initial couplings: $\bar{\varepsilon}_{initial} = 0$ in black, $\bar{\varepsilon}_{initial} = 0.01$ in red and $\bar{\varepsilon}_{initial} = 0.04$ in green. In (c), (d) and (e), we show all clustering coefficients for $\bar{\varepsilon}_{initial} = 0$, $\bar{\varepsilon}_{initial} = 0.01$ and $\bar{\varepsilon}_{initial} = 0.04$, respectively. Panels (f), (g) and (h) show the assortativity for $\bar{\varepsilon}_{initial} = 0$, $\bar{\varepsilon}_{initial} = 0.01$ and $\bar{\varepsilon}_{initial} = 0.04$, respectively.

IV. CONCLUSIONS

In this paper, we studied the effects of plasticity (STP and STDP) on networks of excitatory coupled Hodgkin-Huxley neurons. Neural plasticity is responsible for alterations in the organisation and structure of the brain, and both play an important role in synaptic weights. Besides, STDP has a longer time scale than STP, so it affects differently the structure and function in brain networks.

We started analysing the effect of STP in a pair of neurons that are initially, either uncoupled or bidirectionally coupled. For initially uncoupled neurons, the action of STP and STDP promotes directed connections among neurons with small spike frequency differences, from the faster to the slower spiking neurons. The increase of the recovery-time shortened the interval of frequency differences where connections are formed. When neurons are initially coupled, their frequency difference is smaller and increases the size of the area of directed connections. We found that STP induces uncoupling, depending on the recovery time: the bigger the recovery time, the smaller the interval of frequency difference that allows for the formation of connections.

Next, we build a neural network with an all-to-all topology. Considering only STDP, the coupling matrix exhibits directed connections from neurons with high to neurons with low spike frequencies. We have shown that due to STP, neural networks equipped with STDP

facilitate the formation of synapses among neurons with similar spike frequencies only and that modular neural networks can emerge as a direct result of the combined effect of STP and STDP, a phenomenal structure also depicted by neurophysiological and experimental studies. However, by increasing the STP recovery time, the number of connections decreased and as a consequence, the modules disappeared. That is actually a way to control the modular organization in neural networks. The structure of these modular networks is complex, unlike those in random or small-world networks, resembling more to networks with preferential attachment properties.

In future, we plan to study neural networks with greater diversity in chemical synapses, addressing other STP and STDP rules related to synapses between excitatory-inhibitory and inhibitory-inhibitory neurons⁵⁷. Finally, another interesting aspect of our work would be the introduction of time delay in the synaptic transmission to study how it affects the evolution of the couplings and the modular properties of neural networks.

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